Boundary element and explicit stabilized methods for the cell-by-cell model in electrophysiology

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Contents

Part I: Boundary element method and Cell-by-Cell model

- Crash course on Boundary Element Method,
- Reduction of Cell-by-Cell model to system of ODEs,
- Numerical experiment.

Part II: Parareal Spectral Deferred Correction and Explicit Stabilized methods

- Parareal Spectral Deferred Correction,
- Explicit stabilized methods,
- Application to the monodomain model.







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Boundary Integral Equations

Let $\Omega \subset \mathbb{R}^d$ a domain and its boundary $\Gamma = \partial \Omega$ be as one of:

 $\underline{(2)}$





Let *u* be any solution to

$$-\Delta u = 0$$
 in Ω .

The Green representation formula gives

$$u(x) = \int_{\Gamma} G(x, y) \partial_n u(y) ds_y$$
$$-\int_{\Gamma} \partial_n G(x, y) u(y) ds_y \quad x \in \Omega,$$

with u(y) the Dirichlet and $\partial_y u(y)$ the Neumann data, and G(x, y) is the fundamental solution.

Collocation Boundary Integral Method

Taking the trace yields

$$u = \mathcal{V}\partial_n u - (\mathcal{K} - \frac{1}{2}I)u$$
 on Γ (1)

with

$$\begin{aligned} \mathcal{V}\rho(x) &= \int_{\Gamma} G(x, y)\rho(y) \mathrm{d}s_{y}, & x \in \Gamma, \\ \mathcal{K}\rho(x) &= \int_{\Gamma} \partial_{n} G(x, y)\rho(y) \mathrm{d}s_{y}, & x \in \Gamma \end{aligned}$$

the single and double layer potentials. Rearranging (1):

$$\mathcal{V}\partial_n u = (\mathcal{K} + \frac{1}{2}I)u, \text{ on } \Gamma.$$
 (2)





and impose (2) on x_j only

$$\mathcal{V}\partial_n u(x_j) = (\mathscr{K} + \frac{1}{2}I)u(x_j) \quad \forall j.$$

We represent $u, \partial_n u$ with trigonometric Lagrangian basis $L_j(x)$, with $L_j(x_k) = \delta_{jk}$:

$$\partial_n u = \sum_{j=1}^M \tilde{u}^j L_j, \quad u = \sum_{j=1}^M u^j L_j$$

Collocation Boundary Integral Method

$$\mathcal{V}\partial_n u(x_j) = (\mathcal{K} + \frac{1}{2}I)u(x_j) \quad \forall j$$

$$\partial_n u = \sum_{j=1}^M \tilde{u}^j L_j, \quad u = \sum_{j=1}^M u^j L_j$$

Matrix formulation

$$V\tilde{\mathbf{u}} = (K + \frac{1}{2}I)\mathbf{u}, \qquad \tilde{\mathbf{u}} = P_S \mathbf{u}$$

with

$$P_S = V^{-1}(K + \frac{1}{2}I)$$

the Poincaré-Steklow operator (or Dirichletto-Neumann map).





Model reformulation

Consider a problem with N cells Ω_i , i = 1, ..., N and *unbounded* extracellular matrix Ω_0 with boundary Γ_0 :

$$\begin{split} &-\sigma_i \Delta u_i = 0, & \text{in } \Omega_i \quad i = 0, \dots, N, \\ &-\sigma_i \partial_n u_i = I_t(V_m, z), & \text{on } \Gamma_{i,0} \quad i = 1, \dots, N, \\ &-\sigma_0 \partial_n u_0 = -I_t(V_m, z), & \text{on } \Gamma_0, \\ &u_i - u_0 = V_m, & \text{on } \Gamma_{i,0} \quad i = 1, \dots, N, \\ &\frac{\mathrm{d} z}{\mathrm{d} t} = g(V_m, z), & \text{on } \Gamma_0, \\ &-\sigma_i \partial_n u_i = \kappa(u_i - u_j), & \text{on } \Gamma_{i,j} \quad 1 \le j, i \le N, \end{split}$$

with
$$I_t(V_m, z) = C_m \frac{\mathrm{d} V_m}{\mathrm{d} t} + I_{\mathrm{ion}}(V_m, z).$$



$$\begin{aligned} &-\sigma_i \Delta u_i = 0, & \text{in } \Omega_i \quad i = 0, \dots, N, \\ &\sigma_0 \partial_n u_0 = I_t(V, z), & \text{on } \Gamma_0, \\ & \frac{\mathrm{d} z}{\mathrm{d} t} = g(V, z), & \text{on } \Gamma_0, \\ &\sigma_i \partial_n u_i + \sigma_j \partial_n u_j = 0, & \text{on } \Gamma_{i,j} \quad 0 \le j < i \le n \\ & u_i - u_j = V, & \text{on } \Gamma_{i,j} \quad 0 \le j < i \le n \\ & \sigma_j \partial_n u_j - \sigma_i \partial_n u_i = 2\kappa V, & \text{on } \Gamma_{i,j} \quad 1 \le j < i \le n \end{aligned}$$



Model discretisation

Discretize the skeleton Γ with M points. Ω Every domain's boundary $\Gamma_i = \partial \Omega_i$ has M_i points.

 $\emptyset \qquad \qquad \text{in } \Omega_i \quad i = 0, \dots, N,$

Recall: $\partial_n u \longrightarrow P_S \mathbf{u}, \quad u \longrightarrow \mathbf{u}.$

 $\sigma_0 P_{S,0} \mathbf{u}_0 = I_t(\mathbf{V}, \mathbf{z}), \qquad \text{on } \Gamma_0,$

 $\begin{aligned} \frac{\mathrm{d} \mathbf{z}}{\mathrm{d} t} &= g(\mathbf{V}, \mathbf{z}), & \text{on } \Gamma_0, \\ \sigma_i P_{S,i} \mathbf{u}_i &+ \sigma_j P_{S,j} \mathbf{u}_j = 0, & \text{on } \Gamma_{i,j} & 0 \leq j < i \leq N, \\ \mathbf{u}_i &- \mathbf{u}_j = \mathbf{V}, & \text{on } \Gamma_{i,j} & 0 \leq j < i \leq N, \\ \sigma_j P_{S,j} \mathbf{u}_j &- \sigma_i P_{S,i} \mathbf{u}_i = 2\kappa \mathbf{V}, & \text{on } \Gamma_{i,j} & 1 \leq j < i \leq N, \end{aligned}$

Now, we need to define some restriction $\Gamma \to \Gamma_i$ and extension $\Gamma_i \to \Gamma$ operators.

$$\begin{aligned} & -\sigma_i \Delta u_i = 0, & \text{in } \Omega_i \quad i = 0, \dots, N, \\ & \sigma_0 \partial_n u_0 = I_t(V, z), & \text{on } \Gamma_0, \\ & \frac{\mathrm{d} z}{\mathrm{d} t} = g(V, z), & \text{on } \Gamma_0, \\ & \sigma_i \partial_n u_i + \sigma_j \partial_n u_j = 0, & \text{on } \Gamma_{i,j} \quad 0 \le j < i \le u_i - u_j = V, \\ & u_i - u_j = V, & \text{on } \Gamma_{i,j} \quad 0 \le j < i \le \sigma_j \partial_n u_j - \sigma_i \partial_n u_i = 2\kappa V, & \text{on } \Gamma_{i,j} \quad 1 \le j < i \le v \end{aligned}$$



Global to local operators

The boolean connectivity matrix $A_i : \mathbb{R}^M \to \mathbb{R}^{M_i}$ maps a global vector on Γ to a local vector on Γ_i . The transpose A_i^{T} maps local to global.









Global to local operators with sign change

The signed boolean connectivity matrix $B_i : \mathbb{R}^M \to \mathbb{R}^{M_i}$ maps a global vector on Γ to a local vector on Γ_i . A sign change occurs if the neighbouring domain has higher index.

















Model discretisation

We transpose the equations below, living Ω on $\Gamma_{i,j}$ and Γ_0 , to the global domain Γ .

With A_g the operator from Γ to the gap junctions.

$$\sigma_0 P_{S,0} \mathbf{u}_0 = I_t(A_0 \mathbf{V}, \mathbf{z}) \qquad \in \mathbb{R}^{M_0} = \frac{\mathrm{d} \mathbf{z}}{\mathrm{d} t} = g(A_0 \mathbf{V}, \mathbf{z}) \qquad \in \mathbb{R}^{M_0} = \frac{\mathrm{d} \mathbf{z}}{\mathrm{d} t}$$

$$\sum_{i=0}^{N} \sigma_{i} A_{i}^{\mathsf{T}} P_{S,i} \mathbf{u}_{i} = 0 \qquad \in \mathbb{R}^{M} = \mathbb{R}^{M}$$

$$\sum_{i=0}^{N} B_i^T \mathbf{u}_i = \mathbf{V} \qquad \in \mathbb{R}^M = \mathbf{I}$$

$$\sum_{i=0}^{N} \sigma_{i} A_{g} B_{i}^{\top} P_{S,i} \mathbf{u}_{i} = -2\kappa A_{g} \mathbf{V} \qquad \in \mathbb{R}^{M_{g}} =$$



Reduction to a DAE system

Goal: dispose of \mathbf{u}_i variables.

$$\sigma_0 P_{S,0} \mathbf{N}_0 = I_t (A_0 \mathbf{V}, \mathbf{z}) \qquad \in \mathbb{R}^{M_0} = \Gamma_0$$
$$\frac{\mathrm{d} \, \mathbf{z}}{\mathrm{d} \, t} = g(A_0 \mathbf{V}, \mathbf{z}) \qquad \in \mathbb{R}^{M_0} = \Gamma_0$$
$$\sum_{i=0}^N \sigma_i A_i \mathbf{P}_{S,i} \mathbf{u}_i = 0 \qquad \in \mathbb{R}^M = \Gamma$$
$$\sum_{i=0}^N \mathbf{B}^T \mathbf{u}_i = \mathbf{V} \qquad \in \mathbb{R}^M = \Gamma$$
$$\sum_{i=0}^N \sigma_i A_g B_i^\top P_{S,i} \mathbf{v}_i = -2\kappa A_g \mathbf{V} \qquad \in \mathbb{R}^{M_g} = \Gamma_g$$



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Procedure: Find maps $\psi_i: \Gamma \to \Gamma_i: V \mapsto \sigma_i P_{S,i} \mathbf{u}_i,$

where \mathbf{u}_i satisfies

$$\sum_{i=0}^{N} \sigma_{i} A_{i}^{\mathsf{T}} P_{S,i} \mathbf{u}_{i} = \mathbf{0}, \qquad \sum_{i=0}^{N} B_{i}^{\mathsf{T}} \mathbf{u}_{i} = \mathbf{V}$$

We obtain the DAE:

$$\psi_0(\mathbf{V}) = I_t(A_0\mathbf{V}, \mathbf{z}) \quad \text{on } \Gamma_0,$$
$$\frac{\mathrm{d} \mathbf{z}}{\mathrm{d} t} = g(A_0\mathbf{V}, \mathbf{z}) \quad \text{on } \Gamma_0,$$

$$\sum_{i=0}^{N} A_{g} B_{i}^{\mathsf{T}} \psi_{i}(\mathbf{V}) = -2\kappa A_{g} \mathbf{V} \quad \text{on } \Gamma_{g}$$

Reduction to a DAE system

Theorem: computing ψ_i

The linear maps $\psi_i(\mathbf{V}) = \sigma_i P_{S,i} \mathbf{u}_i$ satisfy $\psi_i(\mathbf{V}) = -B_i\lambda$ with $\lambda \in \mathbb{R}^M$ and $\beta \in \mathbb{R}^N$ solutions to $\begin{pmatrix} F & G \\ G^{\top} & 0 \end{pmatrix} \begin{pmatrix} \lambda \\ \beta \end{pmatrix} = \begin{pmatrix} \mathbf{V} \\ \mathbf{0} \end{pmatrix}.$

Where

$$F = -\sum_{i=0}^{N} \sigma_{i}^{-1} B_{i}^{\top} (P_{S,i}^{+})^{-1} B_{i}, \quad G = (B_{1}^{\top} \mathbf{e}_{1}, \dots$$

 $P_{S,i}^+ = P_{S,i} + \alpha_i \mathbf{e}_i \mathbf{e}_i^{\mathsf{T}}, \quad \mathbf{e}_i = (1, \dots, 1)^{\mathsf{T}} \in \mathbb{R}^{M_i}, \quad \alpha_i > 0.$



The boundary data \mathbf{u}_i can be computed with

$$\mathbf{u}_i = -\sigma_i^{-1} (P_{S,i}^+)^{-1} B_i \lambda + \beta_i \mathbf{e}_i,$$

where β_0 is free.

 $B_N^{\dagger}\mathbf{e}_N$

Checking correctness of $\psi_i(\mathbf{V}) = \sigma_i P_{S,i} \mathbf{u}_i \approx \sigma_i \partial_{\mathbf{n}} u$

Consider two harmonic functions u_0, u_1 satisfying flux continuity.



Define $V = u_i - u_j$, we recover the fluxes and traces as:

•
$$\sigma_i P_{S,i} \mathbf{u}_i = \psi_i(\mathbf{V}),$$

• $\mathbf{u}_i = -\sigma_i^{-1} (P_{S,i}^+)^{-1} B_i \lambda + \beta_i \mathbf{e}_i.$



Given flux and trace, we compute u_0, u_1 inside Ω_i using the Green representation formula.

$$u(x) = \int_{\Gamma} G(x, y) \partial_n u(y) ds_y - \int_{\Gamma} \partial_n G(x, y) u(y) ds_y$$





Checking correctness of $\psi_i(\mathbf{V}) = \sigma_i P_{S,i} \mathbf{u}_i \approx \sigma_i \partial_{\mathbf{n}} u$

Consider three harmonic functions u_0, u_1, u_2 satisfying flux continuity $(u_1 = u_2)$



Define $V = u_i - u_j$, we recover the fluxes and traces as:

•
$$\sigma_i P_{S,i} \mathbf{u}_i = \psi_i(\mathbf{V}),$$

• $\mathbf{u}_i = -\sigma_i^{-1} (P_{S,i}^+)^{-1} B_i \lambda + \beta_i \mathbf{e}_i.$

Given flux and trace, we compute u_0, u_1, u_2 inside Ω_i using the Green representation formula.





Reduction to an ODE system

Recall that we want to solve

$$\begin{split} \psi_0(\mathbf{V}) &= I_t(A_0\mathbf{V}, \mathbf{z}) & \text{on } \Gamma_0, \\ \frac{\mathrm{d}\,\mathbf{z}}{\mathrm{d}\,t} &= g(A_0\mathbf{V}, \mathbf{z}) & \text{on } \Gamma_0, \\ \sum_{i=0}^N A_g B_i^\top \psi_i(\mathbf{V}) &= -2\kappa A_g \mathbf{V} & \text{on } \Gamma_g. \end{split}$$

Using
$$\psi_i(\mathbf{V}) = -B_i \lambda$$
 yields

$$\sum_{i=0}^N A_g B_i^{\mathsf{T}} B_i \lambda = 2\kappa A_g \mathbf{V},$$

$$A_g \lambda = \kappa A_g \mathbf{V}.$$

With this information we can dispose of the equations on $\Gamma_g.$



Multiply first line of

$$\begin{pmatrix} F & G \\ G^{\mathsf{T}} & 0 \end{pmatrix} \begin{pmatrix} \lambda \\ \beta \end{pmatrix} = \begin{pmatrix} \mathbf{V} \\ \mathbf{0} \end{pmatrix}.$$

with A_0 or A_g , use $A_g \lambda = \kappa A_g \mathbf{V}$ and get

$$\begin{pmatrix} F_{00} & F_{0g} & A_0G \\ F_{g0} & F_{gg} - \kappa^{-1}I & A_gG \\ G^{\mathsf{T}}A_0^{\mathsf{T}} & G^{\mathsf{T}}A_g^{\mathsf{T}} & 0 \end{pmatrix} \begin{pmatrix} \lambda_m \\ \lambda_g \\ \beta \end{pmatrix} = \begin{pmatrix} \mathbf{V}_m \\ \mathbf{0} \\ \mathbf{0} \end{pmatrix} .$$

With $\lambda_m = A_0 \lambda, \lambda_g = A_g \lambda, \mathbf{V}_m = A_0 \mathbf{V}$. Thus

$$\psi_0(\mathbf{V}) = -B_0\lambda = A_0\lambda = \lambda_m$$

and $\psi_0(\mathbf{V})$ is replaced with $\psi(\mathbf{V}_m) = \lambda_m$.

Reduction to an ODI

Recall that:
$$I_t(\mathbf{V}_m, \mathbf{z}) = C_m \frac{\mathrm{d} \mathbf{V}_m}{\mathrm{d} t} + I_{\mathrm{ion}}(\mathbf{V}_m, \mathbf{z}).$$

Theorem: the ODE system.

The spatially discretized Cell-

with $\psi(\mathbf{V}_m) = \lambda_m$ solution to

E system

$$I_{ion}(\mathbf{V}_{m}, \mathbf{z}).$$

$$I_{ion}$$



We consider the extracellular matrix and 10 cells:



And solve

$$\psi(\mathbf{V}_m) = C_m \frac{\mathrm{d} \, \mathbf{V}_m}{\mathrm{d} \, t} + I_{\mathrm{ion}}(\mathbf{V}_m, \mathbf{z}) \qquad \text{on } \Gamma_0,$$
$$\frac{\mathrm{d} \, \mathbf{z}}{\mathrm{d} \, t} = g(\mathbf{V}_m, \mathbf{z}) \qquad \text{on } \Gamma_0.$$



With ionic model

$I_{\text{ion}}(V) = \eta_0 V(1 - V/V_{th})(1 - V/V_p),$

without gating variables.

We stimulate the first cell:







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Spectral Deferred Correction method¹

Consider

$$y' = f(y), \qquad y(0) = y_0$$

and an approximation $\tilde{y}(t)$ to the solution y(t). Let

$$\delta(t) = y(t) - \tilde{y}(t)$$

be the error and

$$\varepsilon(t) = y_0 + \int_0^t f(\tilde{y}(s)) ds - \tilde{y}(t)$$

the residual. Then

$$\begin{split} \delta(t_2) &= \delta(t_1) + \int_{t_1}^{t_2} f(\tilde{y}(s) + \delta(s)) - f(\tilde{y}(s)) \mathrm{d}s \\ &+ \varepsilon(t_2) - \varepsilon(t_1) \,. \end{split}$$

¹Dutt, A., Greengard, L., Rokhlin, V. (2000). BIT Numerical Mathematics, 40(2).



Spectral Deferred Correction (SDC) method:

- Fix collocation points c_1, \ldots, c_s in $[t_n, t_n + \Delta t]$,
- Compute approximations \tilde{y}_i at c_i ,

Then iterate on:

- Interpolate and form $\tilde{y}(t) = \sum L_i(t)\tilde{y}_i$,
- Approximate $\varepsilon(t)$ with <u>care</u>,
- Compute δ_i and correct $\tilde{y}_i + \delta_i \rightarrow \tilde{y}_i$.

Hybrid Parareal Spectral Deferred Correction method² U

Parareal SDC:





² Minion, M., Williams, S. 2008, 2010.

Parareal Spectral Deferred Correction method²



² Minion, M., Williams, S. 2008, 2010.



Parareal Spectral Deferred Correction method²

Parareal SDC:









The Second Kind Runge-Kutta-Chebyshev method

One step of RKU is given by

$$\begin{split} k_0 &= y_0, \qquad k_1 = k_0 + \mu_1 \Delta t \; f(k_0), \\ k_j &= \nu_j \; k_{j-1} + \kappa_j \; k_{j-2} + \mu_j \Delta t \; f(k_{j-1}), \quad j = 2, \dots, s, \\ y_1 &= k_s, \end{split}$$

with s satisfying $\Delta t \rho(\partial f/\partial y) \leq (2/3)s(s+2)$.



³ Abdulle, A., Grote M., Rosilho G. 2022. Math. Comput. (in press). ⁴ Croci M., Rosilho G. 2022. J. Comput. Phys. 464. G. Rosilho de Souza



- No step size restriction: just increase *s*.
- Fully explicit,
- There is a multirate version³ for

$$y' = f_F(y) + f_S(y) \, .$$

Good for multiscale ionic models or nonuniform grids, for instance.

- Works in mixed-precision arithmetic⁴ (also in multirate). Good for CPU, memory, and energy savings in HPC computations.
- All flavors are straightforward to implement.





The Parareal SDC RKU method

- Fix collocation points c_1, \ldots, c_s in $[t_n, t_n + \Delta t]$ (Lobatto, Radau,...),
- Compute approximations \tilde{y}_i at c_i with RKU.

Then iterate on:

- Define $\tilde{y}(t) = \sum L_i(t)\tilde{y}_i$,
- Approximate $\varepsilon(t) \approx \sum L_i(t)\varepsilon(c_i)$. $\varepsilon(c_i)$ computed with Lobatto, Radau,... quadrature rules.
- Compute δ_i at c_1, \ldots, c_s solving the error equation with RKU

$$\begin{split} d_0 &= \delta_i, \qquad d_1 \ \overline{\mathcal{T}}(\underline{d}_2) \pm \mu_{\delta}(\underline{A}_1^t \ (f(\tilde{y}^{j-1} + d_{j-1}) - f(\tilde{y}^{j-1})) + \mathcal{E}_f^1(\overline{y}(\underline{s}))) ds \\ d_j &= \nu_j \ d_{j-1} + \kappa_j \ d_{j-2} + \mu_j \Delta t \ (f(\tilde{y}^{j-1} + d_{j-1}) - f(\tilde{y}^{j-1})) + \mathcal{E}_j^j - \nu_j \mathcal{E}^{j-1} - \kappa_j \mathcal{E}^{j-2}, \\ &+ \mathcal{E}(t_2) - \mathcal{E}(t_1) \ . \end{split}$$





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 $\begin{array}{ll} \text{Consider } \Omega = [0,5]cm, \ T = 480ms \ \text{and} \\ \\ \partial_t u = \nu \Delta u - I_{ion}(u,z) + I_s(t), & \text{in } \Omega \times [0,T] \\ \\ z' = g(u,z), & \text{in } \Omega \times [0,T] \end{array}$

With periodic boundary conditions on u, $\nu = 10^{-3}$, I_{ion} , g an ionic model and z its state variables. V





- Discretize with finite differences,
- Solve with Parareal SDC using EE, RKU, and mRKU.

- We use $160 \times 3ms$ subintervals (cores)
- 6 Lobatto collocation nodes on coarse grid,
- 40 Lobatto collocation nodes on fine grid.











<u>Costs per iter per time slice</u> $[t_n, t_n + \Delta t]$

On coarse grid • • •

Cost EE: ≈ 305

Cost RKU: ≈ 97

<u>On fine grid</u> ·····

Cost EE: ≈ 327

Cost RKU: ≈ 190



Multirate RKU method

Consider

$$y' = f_F(y) + f_S(y), \qquad y(0) = y_0,$$

with f_F stiff but cheap and f_S mildly stiff but expensive.

For RKU, the number of costly f_S evaluations is dictated by a few stiff terms in f_F .

We solve the modified problem

$$y'_{\eta} = f_{\eta}(y_{\eta}), \qquad y(0) = y_0,$$

With $\eta \ge 0$ a parameter used to tune the stiffness. For $\eta = \mathcal{O}(\rho_S^{-1})$ and the stiffness of f_{η} is same as f_S .



The averaged force is defined as $f_{\eta}(y) = \frac{1}{\eta} \left(u(\eta) - y \right)$ With auxiliary solution u given by $u' = f_F(u) + f_S(y), \qquad u(0) = y.$

The multirate RKU method is given by:

- Integrate $y'_{\eta} = f_{\eta}(y_{\eta})$ with a RKU method.
- To evaluate f_{η} solve $u' = f_F(u) + f_S(y)$ with another RKU method.



$$f(u,z) = \begin{pmatrix} \nu \Delta u - I_{ion}(u,z) + I_s(t) \\ g(u,z) \end{pmatrix} \quad f_F(u,z) = \begin{pmatrix} \nu \Delta u \\ 0 \end{pmatrix} \quad f_S(u,z) = \begin{pmatrix} -I_{ion}(u,z) + I_s(t) \\ g(u,z) \end{pmatrix}$$











Costs per iter per time slice

F. Eval.	EE	RKU	mRKU
f_S	632	287	122
f_F	632	287	287



Very preliminary results!!!







Instabilities?

0 0 5

Roundoff errors?

Bibliography

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Thank you!

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