

# Multirate explicit stabilized integrators for stiff differential equations

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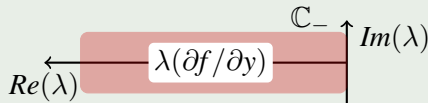
Basel, 1st November 2019.

# Explicit stabilized integrators for stiff differential equations

# Motivating explicit stabilized methods

## Stiff ordinary differential equation

$$y' = f(y), \quad t > 0,$$
$$y(0) = y_0.$$



## Explicit Euler

$$y_{n+1} = y_n + \tau f(y_n)$$

- Straightforward to implement,
- cheap to evaluate.

## Implicit Euler

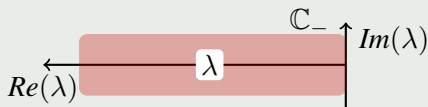
$$y_{n+1} = y_n + \tau f(y_{n+1})$$

- Needs non linear solver routine, preconditioners,
- expensive.

# Motivating explicit stabilized methods

## Dahlquist test equation

$$\begin{aligned}y' &= \lambda y, \quad t > 0, \\ y(0) &= y_0.\end{aligned}$$



## Explicit Euler

$$\begin{aligned}y_{n+1} &= (1 + \tau\lambda)y_n \\ &= R(\tau\lambda)y_n\end{aligned}$$

- $R(z) = 1 + z$ ,
- $|R(z)| \leq 1$  for  $z \in [-2, 0]$ ,
- stability condition  $\tau \leq \frac{2}{|\lambda|}$ .

## Implicit Euler

$$\begin{aligned}y_{n+1} &= (1 - \tau\lambda)^{-1}y_n \\ &= R(\tau\lambda)y_n\end{aligned}$$

- $R(z) = (1 - z)^{-1}$
- $|R(z)| \leq 1$  for all  $\text{Re}(z) \leq 0$ ,
- unconditionally stable.

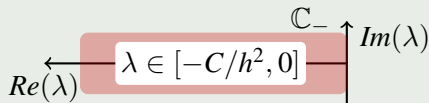
# Motivating explicit stabilized methods

## Space discretized parabolic equation

$$y' = \Delta_h y, \quad t > 0,$$

$$y(0) = y_0,$$

where  $h$  is the smallest element size.



## Explicit Euler

$$y_{n+1} = (I + \tau \Delta_h) y_n,$$

with

$$\tau \leq \frac{2}{|\lambda|} = \mathcal{O}(h^2).$$

## Implicit Euler

$$y_{n+1} = (I - \tau \Delta_h)^{-1} y_n,$$

hence

large system to solve.

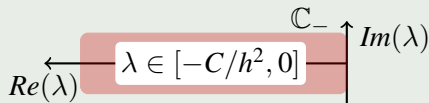
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# Construction of explicit stabilized Runge–Kutta methods

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Given a fixed number of stages  $s$ , find a **first order explicit** scheme with **maximal stability domain along the negative real axis**.

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Given a fixed number of stages  $s$ , find a **first order explicit** scheme with **maximal stability domain along the negative real axis**.

The stability polynomial of such a scheme solves the following

Optimization problem (Markoff, 1916; Guillou and Lago, 1960)

Find a **polynomial**  $R_s(x)$  of degree  $s$  satisfying

$$R_s(0) = R'_s(0) = 1$$

and

$$|R_s(z)| \leq 1 \text{ for } z \in [-\ell_s, 0] \text{ with } \ell_s \text{ maximal.}$$

# Solution to the optimization problem

Chebyshev polynomials of the first kind are defined recursively by

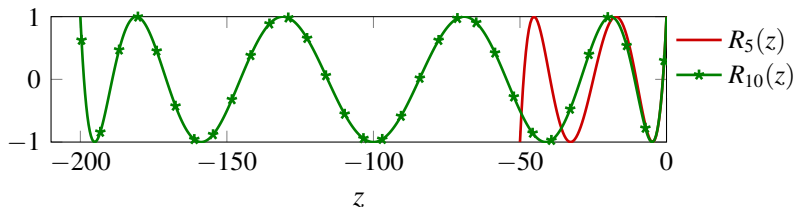
$$T_0(x) = 1, \quad T_1(x) = x, \quad T_n(x) = 2xT_{n-1}(x) - T_{n-2}(x)$$

and satisfy

$$T_n(1) = 1, \quad T'_n(1) = n^2, \quad |T_s(x)| \leq 1 \text{ for all } x \in [-1, 1].$$

Thus,

$$R_s(z) = T_s\left(1 + \frac{z}{s^2}\right) \text{ satisfies } \begin{cases} R_s(0) = R'_s(0) = 1, \\ |R_s(z)| \leq 1 \quad \forall z \in [-2s^2, 0]. \end{cases}$$



# Consequences

- For each  $s$ , there is a first order accurate polynomial  $R_s(z)$  satisfying  $|R_s(z)| \leq 1$  for all  $z \in [-2s^2, 0]$ .
- If there exists a Runge–Kutta scheme having  $R_s(z)$  as stability polynomial, the stability condition on that scheme would be

$$\begin{array}{lll} \tau\lambda \in [-2s^2, 0] & \forall \lambda = \lambda \left( \frac{\partial f}{\partial y} \right) & \Longleftrightarrow \\ \tau\rho \leq 2s^2 & \rho = \rho \left( \frac{\partial f}{\partial y} \right). & \end{array}$$

- If such a scheme exists for all  $s$ , instead of adapting the step size  $\tau$  we can change scheme and take  $s$  larger.
- There is no step size restriction.
- The size of the stability domain of this *family* of Runge–Kutta schemes grows *quadratically* with the number of stages.

# Does such a method exists?

First solution was given by Guillou and Lago (1960), the idea is to write

$$R_s(z) = \prod_{j=1}^s \left(1 - \frac{1}{z_i} z\right), \quad \text{with} \quad z_i \text{ roots of } R_s(z).$$

And represent the scheme as composition of Euler steps:

$$\begin{aligned} k_0 &= y_n, \\ k_j &= k_{j-1} - \frac{1}{z_i} \tau f(k_{j-1}) \quad \text{for } j = 1, \dots, s, \\ y_{n+1} &= k_s. \end{aligned}$$

Disadvantage: when  $|z_i|$  is small we do a large Euler step and the internal stages  $k_j$  become unstable.

Solution by Lebedev (1994): sort the roots, group them two-by-two and use quadratic factors. Becomes tricky to implement.

# Does such a method exists?

A better solution was given by Van der Houwen and Sommeijer (1980), which uses the recursive property

$$T_0(x) = 1, \quad T_1(x) = x, \quad T_n(x) = 2xT_{n-1}(x) - T_{n-2}(x).$$

## Runge–Kutta–Chebyshev (RKC) method

Set  $s \in \mathbb{N}$  such that  $\tau\rho \leq 2s^2$ . Iterate

$$k_0 = y_n,$$

$$k_1 = k_0 + \mu_1 \tau f(k_0),$$

$$k_j = \nu_j k_{j-1} - \kappa_j k_{j-2} + \mu_j \tau f(k_{j-1}) \quad \text{for } j = 2, \dots, s,$$

$$y_{n+1} = k_s.$$

For  $y' = \lambda y$  and  $z = \tau\lambda$  it holds

$$k_j = T_j \left( 1 + \frac{z}{s^2} \right) y_n \quad \text{and thus} \quad y_{n+1} = T_s \left( 1 + \frac{z}{s^2} \right) y_n = R_s(z) y_n.$$

# Cost of the RKC method

We estimate the cost, in the number of function evaluations, when integrating from  $t = 0$  to  $t = 1$ .

- For RKC: take  $\tau = 1$ , since  $\tau\rho \leq 2s^2$  then  $s = \sqrt{\rho/2}$ :

$$C_{RKC} = s = \sqrt{\frac{\rho}{2}}.$$

- For explicit Euler: take  $\tau = 2/\rho$  and  $1/\tau$  time steps:

$$C_{EE} = \frac{1}{\tau} = \frac{\rho}{2}.$$

- For  $\rho = C/h^2$ :

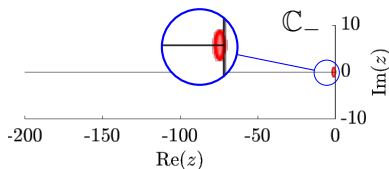
$$C_{RKC} = \sqrt{\frac{C}{2}} \frac{1}{h}, \quad C_{EE} = \frac{C}{2} \frac{1}{h^2}.$$

- Comparison with implicit Euler depends on a multitude of factors: system size, non linearity, preconditioners, parallelism,...

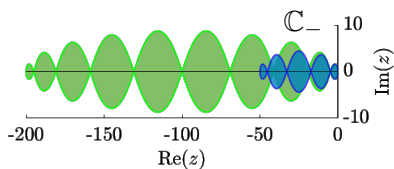
# Stability domain

$$\mathcal{S} = \{z \in \mathbb{C}_- : |R_s(z)| \leq 1\}$$

$\mathcal{S}$  for explicit Euler.



$\mathcal{S}$  for  $s = 10, s = 5$ .

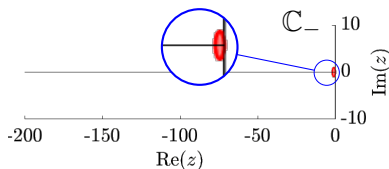


- Problem: unstable in imaginary direction.

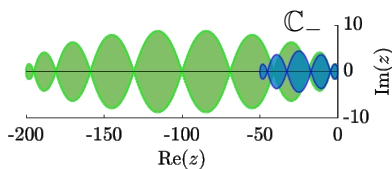
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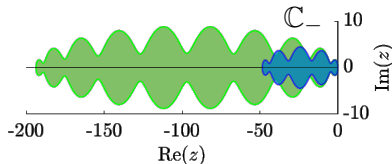


- Problem: unstable in imaginary direction.

- Replace  $T_s \left(1 + \frac{z}{s^2}\right)$  by

$$R_s(z) = \frac{T_s(\omega_0 + \omega_1 z)}{T_s(\omega_0)}$$

$\mathcal{S}$  for  $s = 10, s = 5$ . Damped.



# Numerical experiment

$$\begin{aligned} \text{Solve} \quad & \partial_t u - \Delta u = f && \text{in } \Omega_\delta \times [0, T], \\ & \nabla u \cdot \mathbf{n} = 0 && \text{in } \partial\Omega_\delta \times [0, T], \\ & u = 0 && \text{in } \Omega_\delta \times \{0\}, \end{aligned}$$

in a domain  $\Omega_\delta$  containing a narrow channel of width  $\delta$

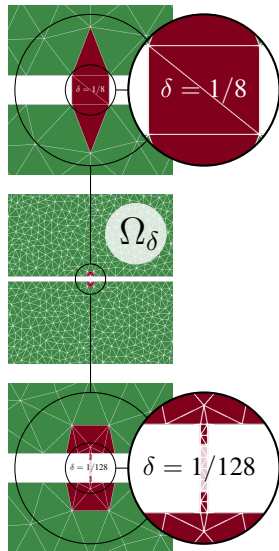
- with first order DG in space

$$\partial_t u_h = \Delta_h u_h + f_h$$

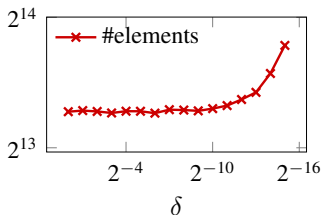
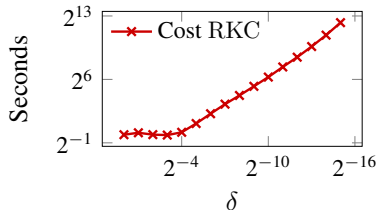
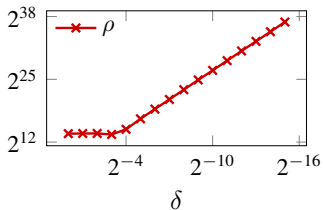
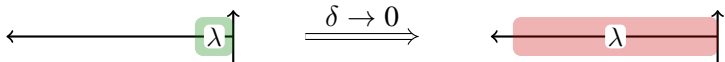
- and RKC in time

$$\tau \rho_h \leq 2s^2 \quad \implies \quad s = \mathcal{O}(h^{-1})$$

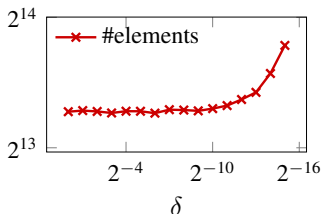
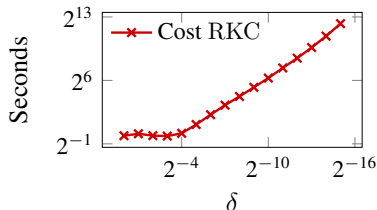
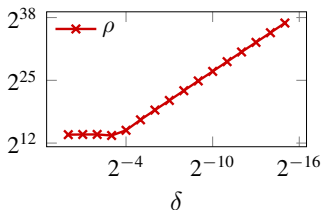
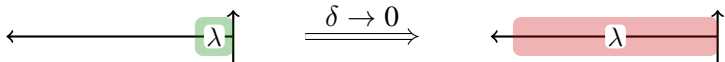
We fix  $\tau = 0.01$  and monitor the performance of RKC as  $\delta \rightarrow 0$ .



# Numerical experiment



# Numerical experiment



## Conclusion

Stabilization of modes induced by a very few degrees of freedom comes at huge computational cost.

# Multirate explicit stabilized methods

# Problem statement

## Multirate equation

Solve the dissipative system

$$y' = f_F(y) + f_S(y), \quad t > 0,$$
$$y(0) = y_0.$$

Term	Stiff ?	Cost ?
$f_F$	stiff	cheap
$f_S$	not stiff	expensive
$f_F + f_S$	stiff	expensive

Examples:

- chemical systems with many slow reactions and a few fast reactions,
- highly integrated electrical circuits with latent and active components,
- parabolic problems on locally refined meshes,
- ...

# Parabolic problem on locally refined mesh

Solve

$$\partial_t u - \Delta u + \beta \cdot \nabla u + \mu u = 0.$$

Space discretization gives:

$$\begin{pmatrix} \partial_t u_h \\ \partial_t u_H \end{pmatrix} = \underbrace{\begin{pmatrix} \Delta_h u_h \\ 0 \end{pmatrix}}_{\text{Fast}} + \underbrace{\begin{pmatrix} \beta \cdot \nabla_h u_h + \mu u_h \\ \Delta_H u_H + \beta \cdot \nabla_H u_H + \mu u_H \end{pmatrix}}_{\text{Slow}}$$

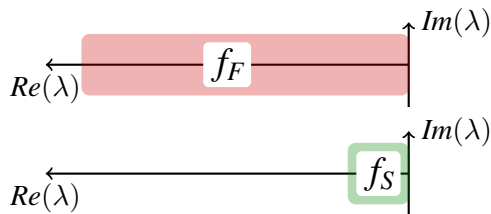
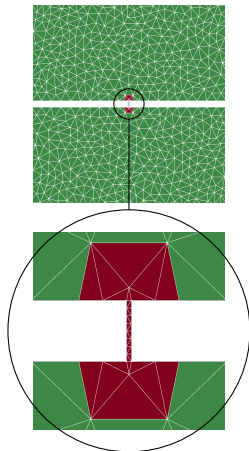


Figure. Spectrum of  $\Delta_h$  and  $\Delta_H$ .



# Problem statement

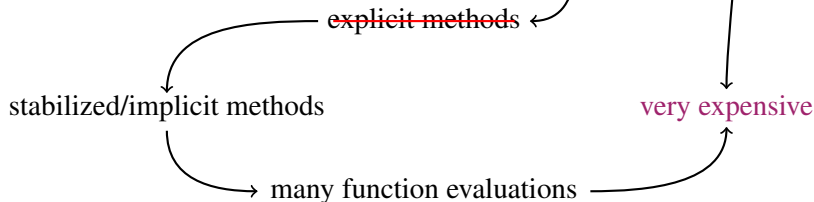
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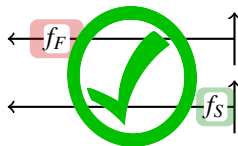
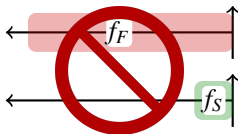
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$f_F$	stiff	cheap
$f_S$	not stiff	expensive
$f_F + f_S$	stiff	expensive

Integration with schemes for  $y' = f(y)$ :



Most of existing multirate methods

- have a spectrum clustering assumption, that is a clear partition between fast and slow variables (E, 2003),

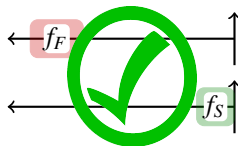
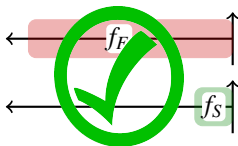


- coupling of fast and slow variables done by *interpolation* or extrapolation  $\Rightarrow$  prone to *instabilities* and/or reduction of stability domain (Gear and Wells, 1984),
- when stable require solution of large and complex non linear systems (Ewing et al., 1990).

# New explicit stabilized multirate method

Multirate RKC<sup>2</sup> method (Abdulle, Grote and Rosilho, 2019):

- no assumption on spectrum clustering,



- no interpolations,
- fully explicit,
- proven to be stable on a large region close to the negative real axis.

# Modified multirate equation

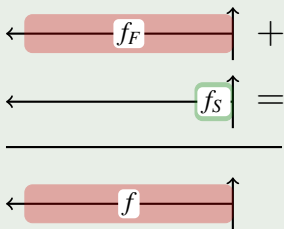
## Idea

Shrink spectrum of  $f_F$  and integrate the modified system.

## Original equation

$$y' = f(y) = f_F(y) + f_S(y).$$

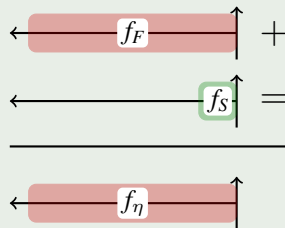
Spectral properties:



## Modified equation

$$y'_\eta = f_\eta(y_\eta) \quad \text{with } \eta \geq 0.$$

For  $\eta = 0$  it holds  $f_\eta = f$  hence:



# Modified multirate equation

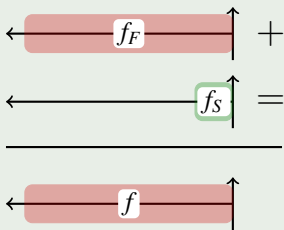
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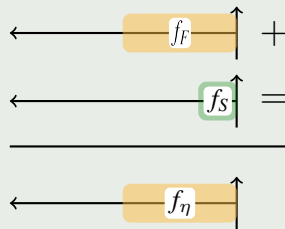
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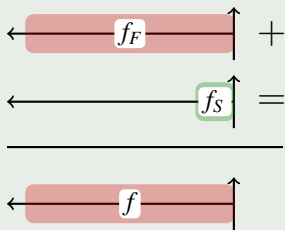
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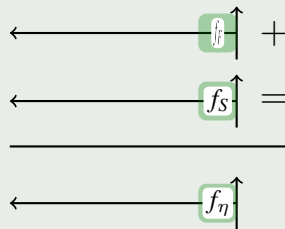
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# Motivating the definition of $f_\eta$


## Properties of $f_\eta$

- $f_\eta = f + \mathcal{O}(\eta),$

- $\rho_\eta \ll \rho.$

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■  $f_\eta = f + \mathcal{O}(\eta)$ , 

■  $\rho_\eta \ll \rho$ .

## Towards the definition of $f_\eta$

Let  $u_0 \in \mathbb{R}^n$  and  $u : [0, \eta] \rightarrow \mathbb{R}^n$  such that


$$u(0) = u_0, \quad \text{and} \quad u \text{ is smooth.}$$


Let

$$f_\eta(u_0) = \frac{1}{\eta} \int_0^\eta f(u(s)) \, ds.$$

# Motivating the definition of $f_\eta$

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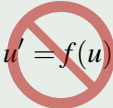
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and


$$u' = f(u).$$



Let

$$f_\eta(u_0) = \frac{1}{\eta} \int_0^\eta f(u(s)) \, ds = \frac{1}{\eta} (u(\eta) - u_0).$$

# Definition of $f_\eta$

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## Definition of $f_\eta$

Let  $u_0 \in \mathbb{R}^n$  and  $u : [0, \eta] \rightarrow \mathbb{R}^n$  such that

$$u(0) = u_0, \quad \text{and} \quad u' = f_F(u) + f_S(u_0).$$

Let

$$f_\eta(u_0) = \frac{1}{\eta} \int_0^\eta f_F(u(s)) \, ds + f_S(u_0) = \frac{1}{\eta} (u(\eta) - u_0).$$

Advantages:

- Computations are cheap since the expensive term  $f_S$  is frozen.
- Stiffness is reduced since  $f_F$  is not frozen.

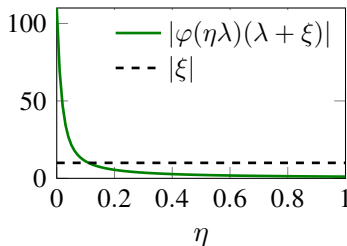
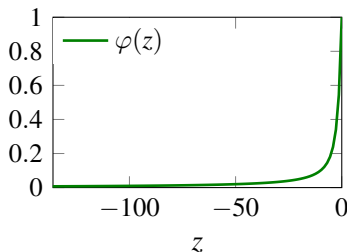
# Stability analysis

Let the multirate Dahlquist equation be defined by

$$y' = f_F(y) + f_S(y) = \lambda y + \xi y, \quad \lambda, \xi \leq 0.$$

Then  $u' = f_F(u) + f_S(u_0) = \lambda u + \xi u_0$  and it holds

$$f_\eta(u_0) = \varphi(\eta\lambda)(\lambda + \xi)u_0, \quad \text{with} \quad \varphi(z) = \frac{e^z - 1}{z}.$$



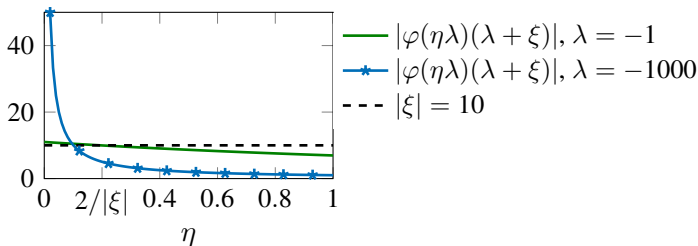
**Goal:** Choose  $\eta$  such that spectrum of  $f_\eta$  is similar to the one of  $f_S$ .  
Hence, we want  $|\varphi(\eta\lambda)(\lambda + \xi)| \leq |\xi|$ .

# Stability analysis

## Lemma

*It holds  $|\varphi(\eta\lambda)(\lambda + \xi)| \leq |\xi|$  for all  $\lambda \leq 0$  if and only if  $\eta \geq 2/|\xi|$ .*

- For  $\eta \geq 2/|\xi|$  the stiffness of  $f_\eta$  depends only on  $f_S$ .
- $\eta$  depends only on  $\xi$ .
- True for all  $\lambda \leq 0$ , so there is no scale separation assumption.
- For a parabolic equation,  $\lambda$  and  $\xi$  represent the eigenvalues of the laplacians  $\Delta_h$  and  $\Delta_H$ . Since  $\Delta_h$  has large and small eigenvalues it is important that the result holds for all  $\lambda \leq 0$ .



# Modified multirate equation

## Modified multirate equation

Solve

$$y'_\eta = f_\eta(y_\eta), \quad t > 0, \quad y_\eta(0) = y_0$$

with

$$f_\eta(u_0) = \frac{1}{\eta}(u(\eta) - u_0),$$

where  $u$  is defined by

$$u' = f_F(u) + f_S(u_0), \quad t \in ]0, \eta], \quad u(0) = u_0, \quad \eta = 2/\rho_S$$

and  $\rho_S$  is the spectral radius of the Jacobian of  $f_S$ .

# Modified multirate equation

## Modified multirate equation

Solve

$$y'_\eta = f_\eta(y_\eta), \quad t > 0, \quad y_\eta(0) = y_0$$

with

**Integrated numerically  $\Rightarrow$  different stability properties**

$$f_\eta(u_0) = \frac{1}{\eta}(u(\eta) - u_0),$$

where  $u$  is defined by

$$u' = f_F(u) + f_S(u_0), \quad t \in ]0, \eta], \quad u(0) = u_0,$$

$$\eta = 2/\rho_S$$

and  $\rho_S$  is the spectral radius of the Jacobian of  $f_S$ .

# Multirate RKC<sup>2</sup> scheme

## Multirate RKC<sup>2</sup> scheme

Let  $\tau > 0$  be the time step, integrate

$$y'_\eta = \bar{f}_\eta(y_\eta), \quad t > 0, \quad y_\eta(0) = y_0,$$

using an RKC scheme with  $m$  stages, where  $\tau \rho_S \leq 2m^2$ .

The right hand side  $\bar{f}_\eta$  is defined by

$$\bar{f}_\eta(u_0) = \frac{1}{\eta}(\bar{u}_\eta - u_0),$$

where  $\bar{u}_\eta$  is an approximation of  $u(\eta)$ , solution of

$$u' = f_F(u) + f_S(u_0), \quad t \in ]0, \eta], \quad u(0) = u_0, \quad \eta = ?,$$

obtained by one step of RKC with  $s$  stages, where  $\eta \rho_F \leq 2s^2$ .

# Stability analysis of numerical $\bar{f}_\eta$

We apply the scheme to the multi rate Dahlquist equation

$$y' = f_F(y) + f_S(y) = \lambda y + \xi y$$

Hence  $u' = \lambda u + \xi u_0$  and  $s \in \mathbb{N}$  is such that  $\eta|\lambda| \leq 2s^2$ . For  $\bar{u}_\eta$ :

$$k_0 = u_0,$$

$$k_1 = k_0 + \mu_1 \eta (\lambda k_0 + \xi u_0),$$

$$k_j = \nu_j k_{j-1} + \kappa_j k_{j-2} + \mu_j \eta (\lambda k_{j-1} + \xi u_0) \quad \text{for } j = 2, \dots, s,$$

$$\bar{u}_\eta = k_s.$$

It can be shown by recursion that

$$\bar{u}_\eta = (R_s(\eta\lambda) + \Phi_s(\eta\lambda)\eta\xi)u_0,$$

with

$$\Phi_s(z) = \sum_{k=1}^s \beta_k U_k(\omega_0 + \omega_1 z)$$

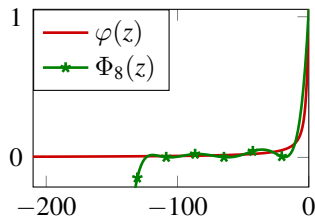
where  $U_k(z)$  is a Chebyshev polynomial of the second kind of degree  $k$  and  $\beta_k, \omega_0, \omega_1$  are parameters of the scheme.

# Stability analysis of numerical $\bar{f}_\eta$

## Lemma

*It holds  $\Phi_s(0) = 1$  and for  $z < 0$*

$$\Phi_s(z) = \frac{R_s(z) - 1}{z}, \quad \varphi(z) = \frac{e^z - 1}{z}$$

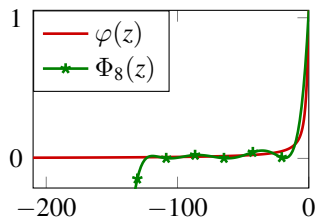


# Stability analysis of numerical $\bar{f}_\eta$

## Lemma

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$$\Phi_s(z) = \frac{R_s(z) - 1}{z}, \quad \varphi(z) = \frac{e^z - 1}{z}$$



Which leads to

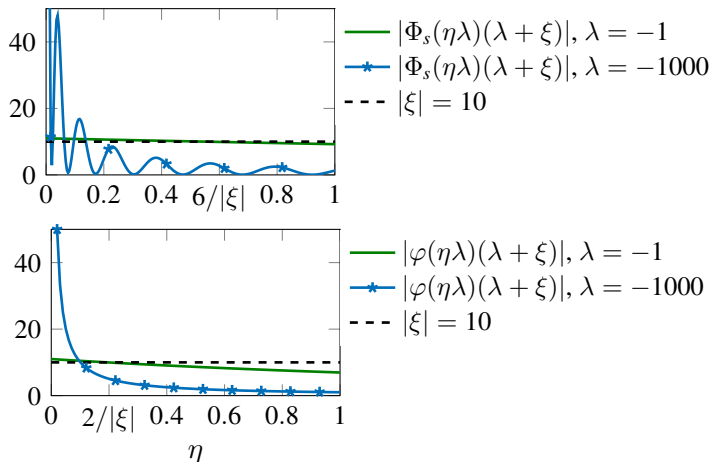
$$\bar{f}_\eta(u_0) = \Phi_s(\eta\lambda)(\lambda + \xi)u_0 \quad f_\eta(u_0) = \varphi(\eta\lambda)(\lambda + \xi)u_0$$

**Goal:** as before, we want the spectrum of  $\bar{f}_\eta$  to be similar to the one of  $f_S$ . Hence, we want  $|\Phi_s(\eta\lambda)(\lambda + \xi)| \leq |\xi|$ .

# Stability analysis of numerical $\bar{f}_\eta$

## Lemma

*It holds  $|\Phi_s(\eta\lambda)(\lambda + \xi)| \leq |\xi|$  for  $\eta\lambda \in [-2s^2, 0]$  if and only if  $\eta \geq 6/|\xi|$ . (we had  $\eta \geq 2/|\xi|$ )*



# Integration of $y' = \bar{f}_\eta$

We apply an RKC scheme to

$$y' = \bar{f}_\eta(y) = \Phi_s(\eta\lambda)(\lambda + \xi)y.$$

Let  $m \in \mathbb{N}$  be such that  $\tau|\xi| \leq 2m^2$ , iterate

$$k_0 = y_n,$$

$$k_1 = k_0 + \mu_1 \tau \Phi_s(\eta\lambda)(\lambda + \xi)k_0,$$

$$k_j = \nu_j k_{j-1} + \kappa_j k_{j-2} + \mu_j \tau \Phi_s(\eta\lambda)(\lambda + \xi)k_{j-1} \quad \text{for } j = 2, \dots, m,$$

$$y_{n+1} = k_m.$$

Hence,

$$y_{n+1} = R_m(\tau \Phi_s(\eta\lambda)(\lambda + \xi))y_n.$$

Since  $|\Phi_s(\eta\lambda)(\lambda + \xi)| \leq |\xi|$  then  $|R_m(\tau \Phi_s(\eta\lambda)(\lambda + \xi))| \leq 1$ .

# Weakening the condition $\eta \geq 6/|\xi|$

- **Problem:** For  $|\xi| \rightarrow 0$  then  $\eta \rightarrow \infty$ .
- But,

$$|R_m(\tau\Phi_s(\eta\lambda)(\lambda + \xi))| \leq 1$$

already for

$$\tau|\Phi_s(\eta\lambda)(\lambda + \xi)| \leq 2m^2,$$

$|\Phi_s(\eta\lambda)(\lambda + \xi)| \leq |\xi|$  is too strong.

- $\eta \geq 3\tau/m^2$  is enough for stability.
- The parameter  $\eta$  depends on the stabilization procedure for  $|\xi|$ , not on  $|\xi|$  itself.
- In practice

$$\eta = 3\frac{\tau}{m^2} = \mathcal{O}(\tau) \quad \text{and generally} \quad \eta \ll \tau.$$

- It follows that the multirate scheme is first order accurate.

# Multirate RKC<sup>2</sup> scheme

## Multirate RKC<sup>2</sup> scheme

Let  $\tau > 0$  be the time step, integrate

$$y'_\eta = \bar{f}_\eta(y_\eta), \quad t > 0, \quad y_\eta(0) = y_0$$

using an RKC scheme with  $m$  stages, where  $\tau \rho_S \leq 2m^2$ .

The right hand side  $\bar{f}_\eta$  is defined by

$$\bar{f}_\eta(u_0) = \frac{1}{\eta}(\bar{u}_\eta - u_0),$$

where  $\bar{u}_\eta$  is an approximation of  $u(\eta)$ , solution of

$$u' = f_F(u) + f_S(u_0), \quad t \in ]0, \eta] \quad u(0) = u_0, \quad \eta = 3\tau/m^2,$$

obtained by one step of RKC with  $s$  stages, where  $\eta \rho_F \leq 2s^2$ .

## Theorem

*The multirate RKC<sup>2</sup> scheme has first order of accuracy and is stable.*

# Numerical experiment

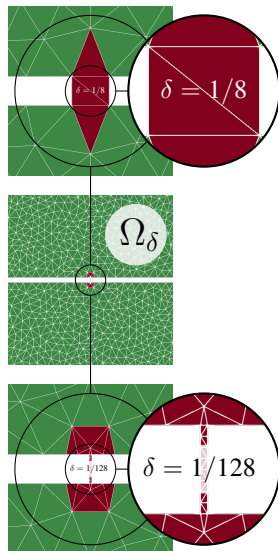
Solve

$$\begin{aligned} \partial_t u - \Delta u &= f && \text{in } \Omega_\delta \times [0, T], \\ \nabla u \cdot \mathbf{n} &= 0 && \text{in } \partial\Omega_\delta \times [0, T], \\ u &= 0 && \text{in } \Omega_\delta \times \{0\}, \end{aligned}$$

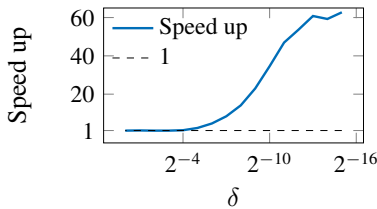
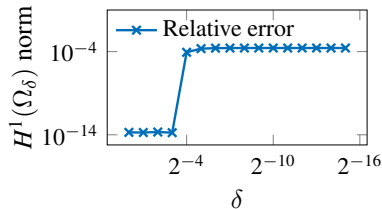
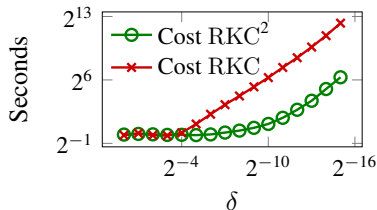
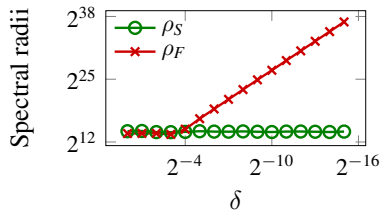
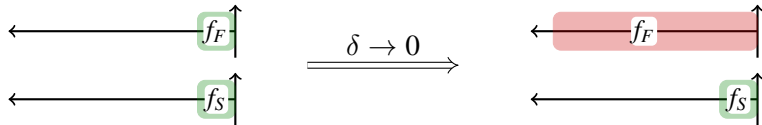
- first order DG in space,
- RKC<sup>2</sup> and RKC in time.

We let  $\delta \rightarrow 0$  and compare the performance of RKC<sup>2</sup> against the one of standard RKC.

$$\begin{pmatrix} \partial_t u_h \\ \partial_t u_H \end{pmatrix} = \underbrace{\begin{pmatrix} \Delta_h u_h \\ 0 \end{pmatrix}}_{\text{Fast}} + \underbrace{\begin{pmatrix} f_h \\ \Delta_H u_H + f_H \end{pmatrix}}_{\text{Slow}}$$



# Numerical experiment



## Thank you for your attention!

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